

## 10.5 Energetics of nuclear reactions

During a nuclear reaction, energy is either evolved or absorbed. Reactions in which energy is evolved are known as *exoergic reactions* while those requiring absorption of energy are called *endoergic*. The total amount of energy evolved or absorbed during a nuclear reaction is called the  $Q$  value or simply the  $Q$  of the reaction. So by definition

$$Q = E_Y + E_y - E_X - E_x = E_Y + E_y - E_x \quad \dots(10.5-1)$$

if the target nucleus  $X$  is at rest.

$Q$  is thus equal to the net surplus (or deficit) of the energies of the reaction products  $E_Y + E_y$  over the energy supplied ( $E_x$ ).

If the atomic masses are expressed in energy units, Eq. (10.3-3) can be rewritten as

$$M_X + M_x + E_x = M_Y + M_y + E_Y + E_y$$

Then we get from Eq. (10.5-1)

$$Q = M_X + M_x - M_Y - M_y \quad \dots(10.5-2)$$

Written in terms of the binding energies of different nuclei, we can also write

$$Q = B_T + B_Y - B_X - B_Z \quad \dots(10.5-3)$$

By definition  $Q > 0$  for an exoergic reaction, while  $Q < 0$  for an endoergic reaction. Since there is a net deficit of energy in the latter case some energy must be supplied for the reaction to occur. This usually comes from the kinetic energy  $E_x$  of the projectile.

Eq. (10.5-2) shows that for an exoergic reaction  $M_T + M_Y$  is greater than  $M_X + M_Z$ , while for an endoergic reaction  $M_T + M_Y$  is less than  $M_X + M_Z$ .

*Threshold energy of an endoergic reaction :*

Eq. (10.5-1) shows that the  $Q$  value of a reaction can be expressed in terms of the kinetic energies of the projectile ( $E_x$ ) and of the product nuclei  $E_T$  and  $E_Y$ .

In view of the energy and momentum conservation laws,  $E_x$  can be expressed in terms of  $E_T$  and  $E_Y$ . Referring to Fig. 10.5, we get from the law of conservation of momentum along and perpendicular to the direction of motion of the projectile ( $\because p = \sqrt{2ME}$ )

$$\sqrt{2M_x E_x} = \sqrt{2M_T E_T} \cos \theta + \sqrt{2M_Y E_Y} \cos \phi \quad \dots(10.5-4)$$

$$0 = \sqrt{2M_x E_x} \sin \theta - \sqrt{2M_T E_T} \sin \phi \quad \dots(10.5-5)$$

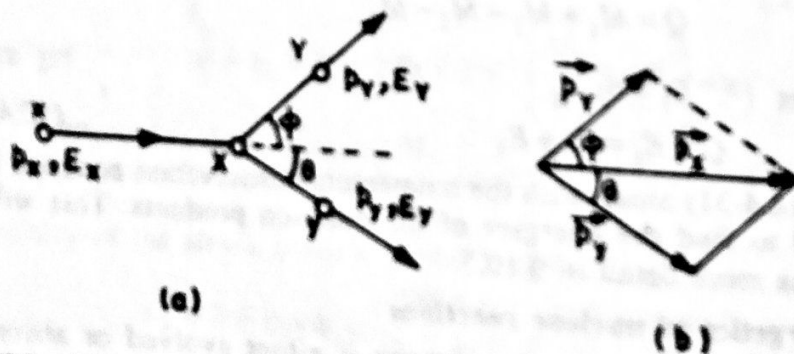


Fig. 10.5. (a) Motion of the projectile (x) and product particles (y and Y) taking part in a nuclear reaction. (b) Momentum diagram.

Eq. (10.5-1) gives the law of conservation of energy :

$$Q = E_T + E_Y - E_x$$

Squaring and adding Eqs. (10.5-4) and (10.5-5), we get

$$2M_T E_T = 2M_x E_x + 2M_Y E_Y - 4 \sqrt{M_x M_Y E_x E_Y} \cos \theta$$

or

$$E_T = \frac{M_x}{M_T} E_x + \frac{M_Y}{M_T} E_Y - \frac{2}{M_T} \sqrt{M_x M_Y E_x E_Y} \cos \theta \quad \dots(10.5-6)$$

Then from Eq. (10.5-1) and (10.5-6) we get

$$Q = E_Y \left( 1 + \frac{M_Y}{M_T} \right) - E_x \left( 1 - \frac{M_x}{M_T} \right) - \frac{2}{M_T} \sqrt{M_x M_Y E_x E_Y} \cos \theta \quad \dots(10.5-7)$$

Eq. (10.5-7) is quadratic in  $z = \sqrt{E_y}$  so that we can write

$$az^2 + bz + c = 0 \quad \dots(10.5-8)$$

where  $a = 1 + \frac{M_y}{M_Y}$ ,  $b = -(2/M_Y) \sqrt{M_x M_y E_x} \cos \theta$

and  $c = -E_x(1 - M_x/M_Y) - Q$

Eq. (10.5-8) has the solution  $z = \frac{-b \pm \sqrt{b^2 - 4ac}}{2a} \quad \dots(10.5-9)$

Written explicitly we then get

$$\sqrt{E_y} = \frac{1}{M_Y + M_y} \left\{ (M_x M_y E_x)^{1/2} \cos \theta \right\} \pm \left[ M_x M_y E_x \cos^2 \theta + (M_Y + M_y) \left\{ Q M_Y + E_x (M_Y - M_y) \right\} \right]^{1/2} \quad \dots(10.5-10)$$

If we write  $Q' = -Q$  then for endoergic reactions,  $Q' > 0$  since  $Q < 0$ . In this case if  $E_x = 0$ , we have

$$b = 0 \quad \text{and} \quad c = -Q = Q' > 0$$

The solution for  $z$  in this case becomes

$$z = \sqrt{E_y} = \pm \frac{\sqrt{-4ac}}{2a} = \pm \sqrt{-Q'/a}$$

Since both  $a$  and  $Q'$  are positive  $z = \sqrt{E_y}$  is imaginary in this case. This means that the reaction is not possible with  $E_x = 0$ . A minimum energy  $E_x = E_{\min}$  is needed to initiate endoergic reaction. In this case the term under the square root sign in Eq. (10.5-9) must be zero so that we get

$$b^2 - 4ac = 0$$

Substituting for  $a$ ,  $b$  and  $c$ , we get

$$\frac{4}{M_Y^2} (M_x M_y E_{\min}) \cos^2 \theta = 4 \left( 1 + \frac{M_y}{M_Y} \right) \left\{ -Q - E_{\min} \left( 1 - \frac{M_x}{M_Y} \right) \right\}$$

which gives

$$E_{\min} = - \frac{(M_y + M_Y) Q}{M_y + M_Y - M_x - (M_x M_y / M_Y) \sin^2 \theta} \quad \dots(10.5-11)$$

Since  $Q < 0$ ,  $E_{\min} > 0$ .

Using Eq. (10.5-2) we get

$$E_{\min} = \frac{-(M_y + M_Y) Q}{M_x - Q - (M_x M_y / M_Y) \sin^2 \theta} \quad \dots(10.5-12)$$

$E_{\min}$  depends on the angle at which the particle  $y$  is emitted. When  $\theta = 0$ , i.e.,  $y$  is emitted in the forward direction,  $E_{\min}$  has the lowest value and is known as the *threshold energy* for the endoergic reaction and is usually written as  $E_{\text{th}}$ . From Eq. (10.5-12) we get

$$E_{\text{th}} = - \frac{(M_y + M_Y) Q}{M_x - Q} \quad \dots(10.5-13)$$



Since  $Q \ll M_x$ , we can neglect it in the denominator of Eq. (11.5-13). Also we can replace  $M_y + M_Y$  in the numerator by  $M_x + M_X$ . So we get finally

$$E_{th} \approx -Q \cdot \frac{M_x + M_X}{M_X} = -Q \left( 1 + \frac{M_x}{M_X} \right) \dots(10.5-14)$$

So by measuring the minimum energy  $E_{th}$  at which an exoergic reaction is initiated it is possible to determine the  $Q$  value of the reaction.

An inspection of Eq. (10.5-10) shows that under certain circumstances  $E_y$  will be a double-valued function of the projectile energy  $E_x$  i.e., for a given  $E_x$ , there may be two values of  $E_y$ , the energy of the emitted particle. This happens only for endoergic reactions. The double valued nature of  $E_y$  is revealed in Fig. 10.6 for the  ${}^3\text{H}(p, n) {}^3\text{He}$  endoergic reaction which has  $Q = -0.7638 \text{ MeV}$ . Eq. (10.5-10) also shows that  $E_y$  is single valued if the following condition is satisfied :

$$QM_Y + E_x (M_Y - M_x) \geq 0$$

or,

$$E_x \geq \frac{-QM_Y}{M_Y - M_x}$$

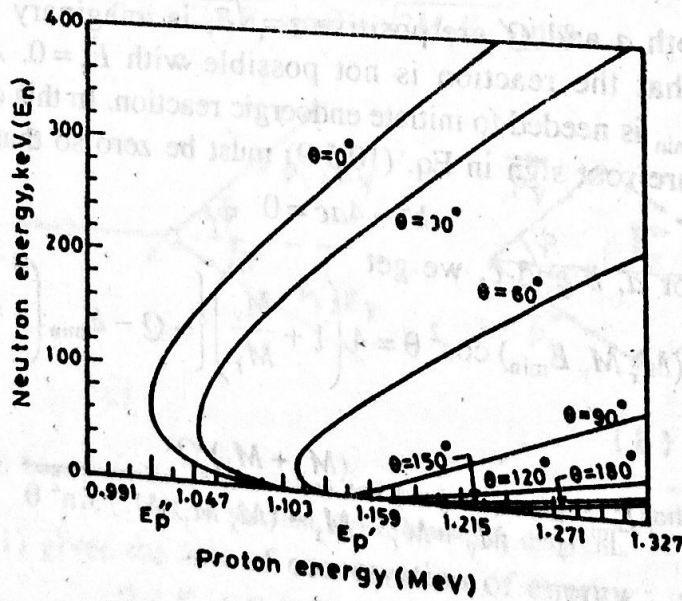


Fig. 10.6.  $E_n$  versus  $E_p$  graph in  ${}^3\text{H}(p, n) {}^3\text{He}$  reaction. Double valued nature of neutron energy should be noted.

Thus there is a limiting energy of the projectile above which the emitted particle energy will be single valued. This is given by

$$E'_x = -\frac{QM_Y}{M_Y - M_x} \dots(10.5-15)$$

For the case cited above  $E'_x = 1.145 \text{ MeV}$ . For projectile energy greater than  $E'_x$ , the product particle  $y$  can be emitted at all angles between  $0^\circ$  and

a maximum angle  $\theta_{\max}$ , which can be found with the help of Eq. (10.5-10).

*Exoergic reaction :*

In this case, the reaction can occur for all values of  $E_x$  including  $E_x = 0$ . For  $E_x = 0$ , the incident momentum is zero and hence the sum of the momenta of the product particles must be zero :  $p_Y + p_y = 0$ . This means that  $Y$  and  $y$  proceed in opposite directions, so that  $\theta + \varphi = \pi$ . Also in this case  $Q = E_y + E_Y$ .

In general Eq. (10.5-7) gives for  $Q > 0$  only one value of  $E_y$ , the energy of the emitted particle, for a given  $E_x$  and at a given angle of emission  $\theta$ . All values of  $\theta$  are possible. Hence there is an energy distribution of the emitted particles between a maximum at  $\theta = 0$  and a minimum at  $\theta = \pi$ . In the solution for Eq. (10.5-7) given by Eq. (10.5-9) the plus sign has to be chosen if a positive value of the momentum  $p_y$  is to be obtained.

As we shall see in § 13.4 *c* exploitation of the rigid correlation between  $E_y$  and  $\theta$  is the only way of obtaining monoenergetic neutron beams of different energies.

## 10.6 Experimental